

Border-lines

26-02-28

$T_{mo} := 100$ $K_r := 10$ $K_{rs} := 5$ $gamL := 2$ $p := 0.01..300$ $v := .5, .501..1$ $x := .5, .501..1$

1. Melting: Lindemann-Form [Li1910]

[Li1910] F. A. Lindemann, Physik. Z.11, 609 (1910)

$$T_{mLv}(v) := T_{mo} \cdot v^{gamL}$$

with the Murnaghan-EOS
[Li1910] gives
with the parametres

$$pM(v) := \frac{K_r}{K_{rs}} \cdot (v^{-K_{rs}} - 1)$$

$$vM(p) := \left(1 + p \cdot \frac{K_{rs}}{K_r}\right)^{-K_{rs}^{-1}}$$

$LSG := gamL \cdot K_{rs}^{-1}$ and $rpM := K_{rs} \cdot K_r^{-1}$

The modified Lindemann -Form;

$$T_{mLi}(p, gamL) := T_{mo} \cdot (1 + rpM \cdot p)^{gamL \cdot K_{rs}^{-1}}$$

The Simon-Glatzel -Form [SG1929]
uses rpM with the free parameter LSG

$$T_{mSG}(p, LSG) := T_{mo} \cdot (1 + rpM \cdot p)^{LSG}$$

Simon F., Glatzel G.. Z.Anorg. Allg. Chem. 178, 309 (1929)

The less restricted Kechin-form [K19951]

uses two free parameters: rpK and LK

$$T_{mKe}(p, T_{mo}, rpK, LK) := T_{mo} \cdot (1 + rpK \cdot p)^{LK}$$

V. V. Kechin, J. Phys.: Condens. Matter 7, 531 (19951)

V. V. Kechin, Phys. Rev. B 65, 052102 (2001)

One reference point: (p_{mr}, T_{mr}) $p_{mr} := 100$ $T_{mr} := 400$ gives for [SG1929]

$LSG := (\log(T_{mr}) - \log(T_{mo})) \cdot (\log(1 + rpM \cdot p_{mr}))^{-1}$ and $gamL := LSG \cdot K_{rs} = 1.763$

and for [K19951]: $LrK(rpK)$

$$LrK(rpK) := (\log(T_{mr}) - \log(T_{mo})) \cdot (\log(1 + rpK \cdot p_{mr}))^{-1}$$

with the free parameter rpK :

$rpK := .5$

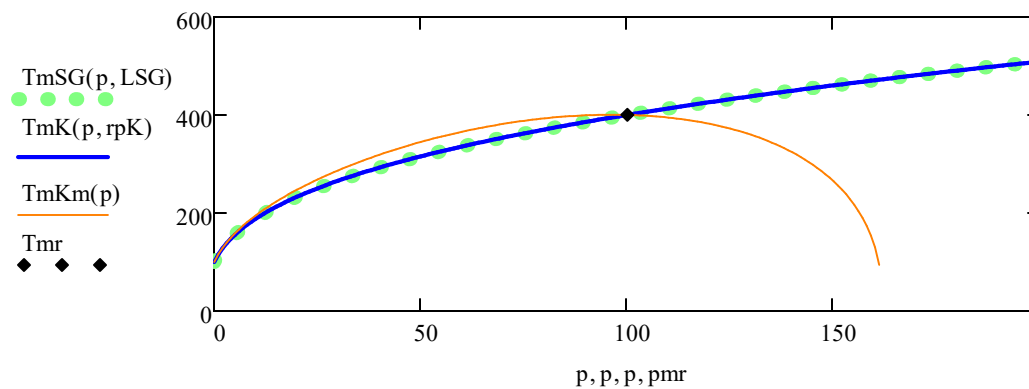
$$T_{mK}(p, rpK) := T_{mo} \cdot (1 + rpK \cdot p)^{LrK(rpK)}$$

For melting with softening [K19951] can be modified with parameters ds and qs :

$ds := .0001$

$qs := .4$

$$T_{mKm}(p) := T_{mK}(p, rpK) \cdot [1 + ds \cdot p \cdot (p_{mr} - p)]^{qs}$$



With the EOS for the phases 1 and 2 or $i = 1..2$
 $V_i(p,T) = V_i(p,0) \times (1 + p_{\theta}(V_i(p,T),T) / K_i(p,T))$

$p_{\theta}(V,T) = (\gamma(V)/V) \times U_{\theta}(V,T)$ with scaled Debye function $u_D(T,\theta(V))$
 $U_{\theta}(V,T) = 3R \times \theta(V) \times u_D(T,\theta(V))$ and
 $V_i(p,T) = V_i(p,0) + dp_{tr}(T) / dT = \Delta S(p_{tr}(T),T) / \Delta V(p_{tr}(T),T)$
 or roughly with V_i

$V_i(p,T) = V_i(p,0) + [3R \times \theta(V_i) \times \gamma(V_i) / K_i(V_i)] \times u_D(T,\theta(V_i))$
 this gives as good approximation for the Volume-difference with
 $[...] = \Delta V_D$ and an average Θ_D :

$$\Delta V_{tr}(T) = \Delta V_{tr}(0) + \Delta V_D \times u_D(T,\Theta_D)$$

and similarly (with $\Delta S_{tr}(0) = 0$)

$$\Delta S_{tr}(T) = \Delta S_D \times s_D(T,\Theta_D) \text{ and finally with the 3 roughly constant parameters}$$

Θ_D , $\Delta S_{V0} = \Delta S_D / V_0 \Delta$, $\Delta V_{D0} = \Delta V_D / V_0$ and the two functions $s_D(T,\Theta_D)$ and $u_D(T,\Theta_D)$:

$$p_{tr}(T) := \Delta S_{V0} \cdot \int_0^{\Theta_D} \frac{s_D(T, \Theta_D)}{1 + \Delta V_{D0} \cdot u_D(T, \Theta_D)} dT$$

Test of best form for $p_{tr}(T)$:

With the Einstein-Forms for $k_B := 1$ $T_E := 100$ $T_D := 1.4 \cdot T_E$ $T := 1..100$

$$f_E(T, T_E) := 3 \cdot k_B \cdot T \cdot \ln\left(1 - \exp\left(\frac{-T_E}{T}\right)\right) \quad f_E(T, T_E) := T \cdot \ln\left(1 - \exp\left(\frac{-T_E}{T}\right)\right)$$

$$s_E(T, T_E) := 3 \cdot k_B \cdot \left[\frac{T_E}{T} \cdot \exp\left(\frac{-T_E}{T}\right) \cdot \left(1 - \exp\left(\frac{-T_E}{T}\right)\right)^{-1} - \ln\left(1 - \exp\left(\frac{-T_E}{T}\right)\right) \right]$$

$$u_E(T, T_E) := 3 \cdot k_B \cdot T_E \cdot \left(\exp\left(\frac{T_E}{T}\right) - 1\right)^{-1} \quad c_E(T, T_E) := 3 \cdot k_B \cdot \left(\frac{T_E}{T}\right)^2 \cdot \exp\left(\frac{T_E}{T}\right) \cdot \left(\exp\left(\frac{T_E}{T}\right) - 1\right)^{-2}$$

one obtains the Debye-Forms for the border-lines:

$$FD_n(T, T_D) := 9 \cdot k_B \cdot \frac{T}{T_D^3} \cdot \int_1^{T_D} T_E^2 \cdot \ln\left(1 - \exp\left(\frac{-T_E}{T}\right)\right) dT_E$$

$$FD(T, T_D) := 9 \cdot k_B \cdot T \cdot \left(\frac{T}{T_D}\right)^3 \cdot \int_0^{\frac{T_D}{T}} z^2 \cdot \ln(1 - \exp(-z)) dz$$

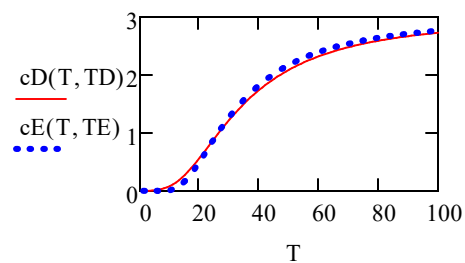
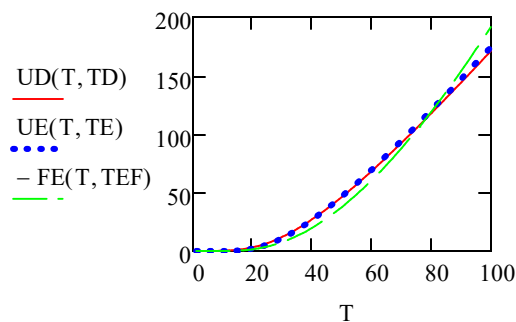
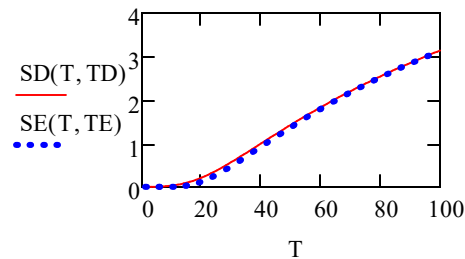
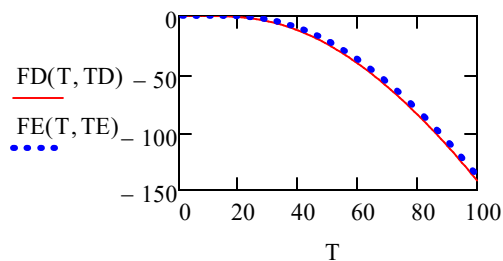
$$SD(T, TD) := 9 \cdot kB \cdot \left(\frac{T}{TD}\right)^3 \cdot \int_0^{\frac{TD}{T}} z^2 \cdot \left[\frac{z}{(\exp(z) - 1)} - \ln(1 - \exp(-z)) \right] dz$$

$$UDn(T, TD) := 9 \cdot kB \cdot TD \cdot \frac{1}{TD^4} \cdot \int_1^{TD} TE^3 \cdot \left(\exp\left(\frac{TE}{T}\right) - 1 \right)^{-1} dTE$$

$$UD(T, TD) := 9 \cdot kB \cdot T \cdot \left(\frac{T}{TD}\right)^3 \cdot \int_0^{\frac{TD}{T}} z^3 \cdot (\exp(z) - 1)^{-1} dz$$

$$cDn(T, TD) := 9 \cdot kB \cdot TD \cdot \frac{1}{T^2 \cdot TD^4} \cdot \int_1^{TD} TE^4 \cdot \exp\left(\frac{TE}{T}\right) \cdot \left(\exp\left(\frac{TE}{T}\right) - 1 \right)^{-2} dTE$$

$$cD(T, TD) := 9 \cdot kB \cdot \left(\frac{T}{TD}\right)^3 \cdot \int_0^{\frac{TD}{T}} z^4 \cdot \exp(z) \cdot (\exp(z) - 1)^{-2} dz \quad fe := .75 \quad TEF := TE \cdot fe$$



Test of best form for ptr(T) of page 4 indicates that UD(T,TD) can be replaced by -FE(TEF) and SD(T,TD) by SE(T,TEF) to give the following forms:

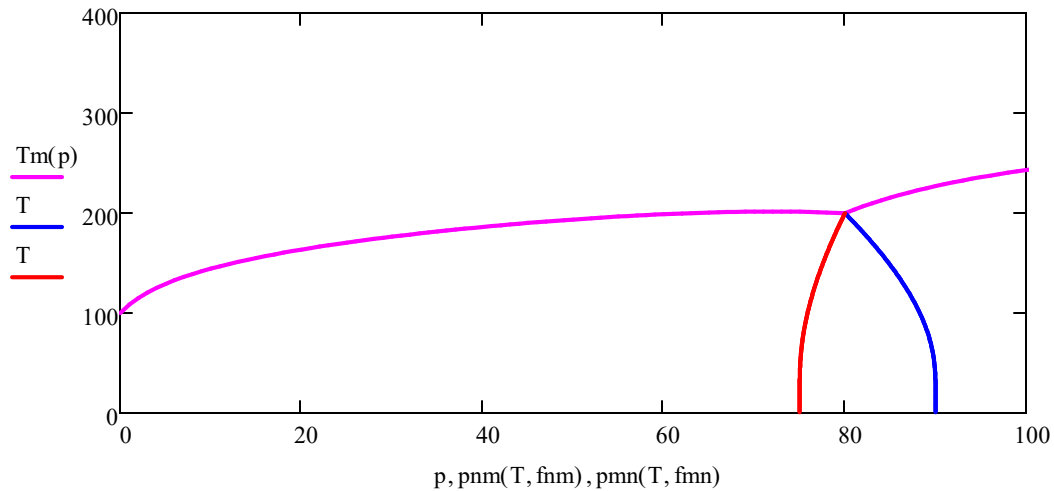
Limiting cases with almost constant : $\Delta V(p, T) \quad T := 1, 2 \dots Tmtri \quad pmtr := ptr = 80$

2a: ptr(T) for two cases (f_{mn}<0, f_{mn}>0) meeting the triple point with pnmo, T_{mo} at the melting curve T_m(p) based on the following data:

$pnmo := 90$ $fnn := .01$ $Tnm := 180$ $pnmo := 75$ $fnn := -.01$ $Tmtri = 200$

$pnm(T, fnn) := pnmo + fnn \cdot FE(T, Tnm)$ $fnn := \text{wurzel}(pnm(Tmtri, fnn) - ptri, fnn) = 0.032$

$pnm(T, fnn) := pnmo + fnn \cdot FE(T, Tnm)$ $fnn := \text{wurzel}(pnm(Tmtri, fnn) - ptri, fnn) = -0.016$

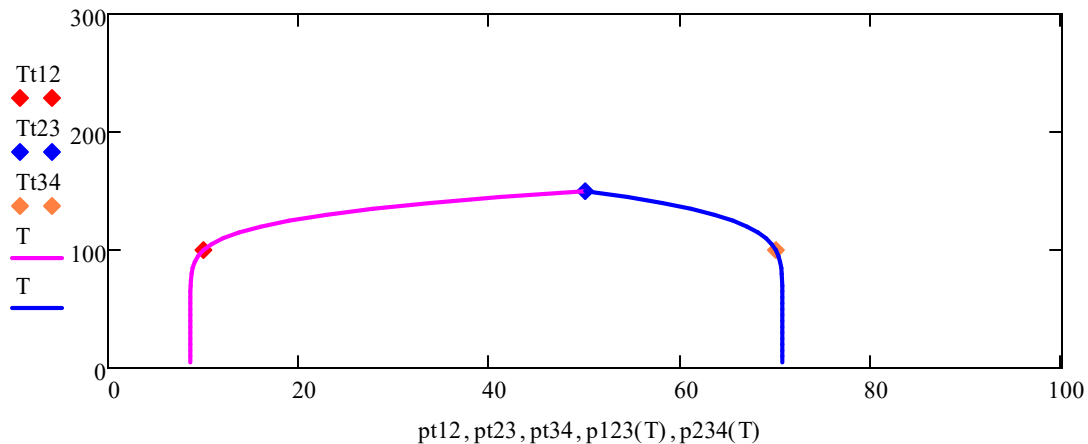


2b: ptr(T) reentrant case based on the following data:

$Tt12 := 100$ $Tt23 := 150$ $pt12 := 10$ $pt23 := 50$ $Tt34 := 100$ $pt34 := 70$ $Tnm := 900$

$p123(T) := pt12 + (pt23 - pt12) \cdot \frac{FE(T, Tnm) - FE(Tt12, Tnm)}{FE(Tt23, Tnm) - FE(Tt12, Tnm)}$ $T := 0, 5.. Tt23$

$p234(T) := pt34 + (pt23 - pt34) \cdot \frac{FE(T, Tnm) - FE(Tt34, Tnm)}{FE(Tt23, Tnm) - FE(Tt34, Tnm)}$



2c: ΔV almost constant, $\Delta F(T)$ based on 1 average (like case 2a) or

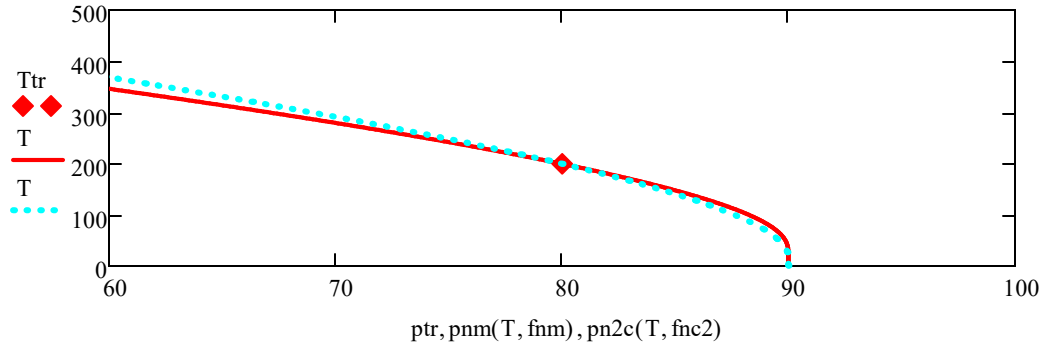
$ptr := ptri$ $Ttr := Tmtri$ $DT := 100$

$fnn = 0.032$ $Tnm := 180$ $pnmo = 90$

**2 individual Einstein functions(2c)
with DT as difference to the single Tnm
give no large difference in the fitted
borderline**

$$\begin{aligned} Tnc1 &:= Tnm + DT & Tnc2 &:= Tnm - DT \\ fnc2 &:= .5 \cdot fnm & T &:= 1, 2 \dots 400 \end{aligned}$$

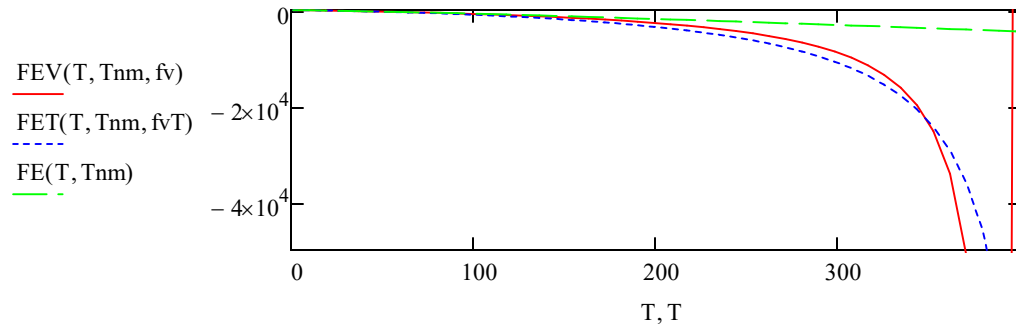
$$\begin{aligned} pn2c(T, fnc2) &:= pnmo + fnc2 \cdot (FE(T, Tnc1) + FE(T, Tnc2)) \\ fnc2 &:= \text{wurzel}(pn2c(Ttr, fnc2) - ptr, fnc2) = 0.012 \end{aligned}$$



3: $\Delta V(T)$ variable by fv for reentrant phases and adjusted FET for ptr(T):

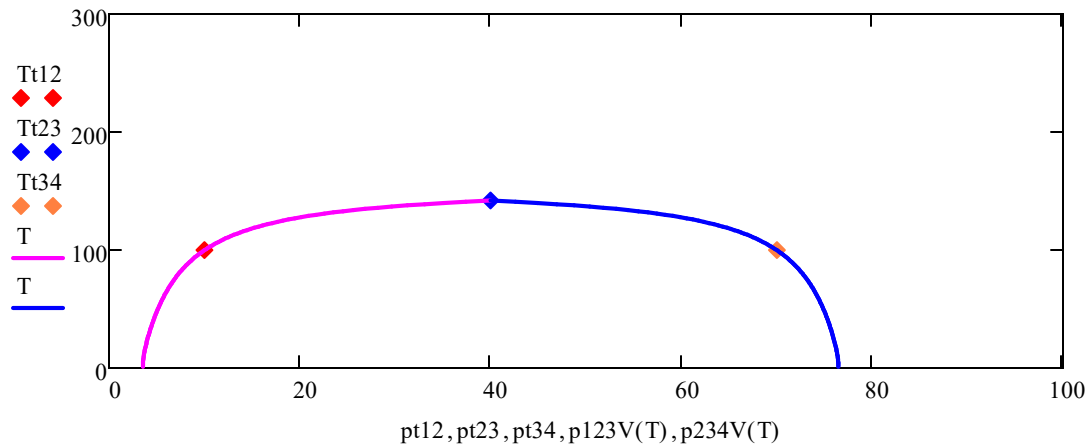
$$\begin{aligned} Tt12 &:= 100 & Tt23 &:= 142 & pt12 &:= 10 & pt23 &:= 40 & Tnm &:= 10 & TnmV &:= 200 & fv &:= .00051 \\ Tt34 &:= 100 & & & pt34 &:= 70 & fv12 &:= .0026 & fv34 &:= .0026 & & & fvT &:= .0024 \\ FEV(T, Tnm, fv) &:= FE(T, Tnm) \cdot (1 - fv \cdot T \cdot SE(T, TnmV))^{-1} & & & & & & & & & & & T &:= 1, 10 \dots 500 \\ FET(T, Tnm, fv) &:= FE(T, Tnm) \cdot (1 - fv \cdot T)^{-1} \end{aligned}$$

**Constant SE (= 1) works well with
adjusted fvT in FET(T,Tnm,fvT) !!!!!**



$$\begin{aligned} p123V(T) &:= pt12 + (pt23 - pt12) \cdot \frac{FEV(T, Tnm, fv12) - FEV(Tt12, Tnm, fv12)}{FEV(Tt23, Tnm, fv12) - FEV(Tt12, Tnm, fv12)} \\ p234V(T) &:= pt34 + (pt23 - pt34) \cdot \frac{FEV(T, Tnm, fv34) - FEV(Tt34, Tnm, fv34)}{FEV(Tt23, Tnm, fv34) - FEV(Tt34, Tnm, fv34)} \end{aligned}$$

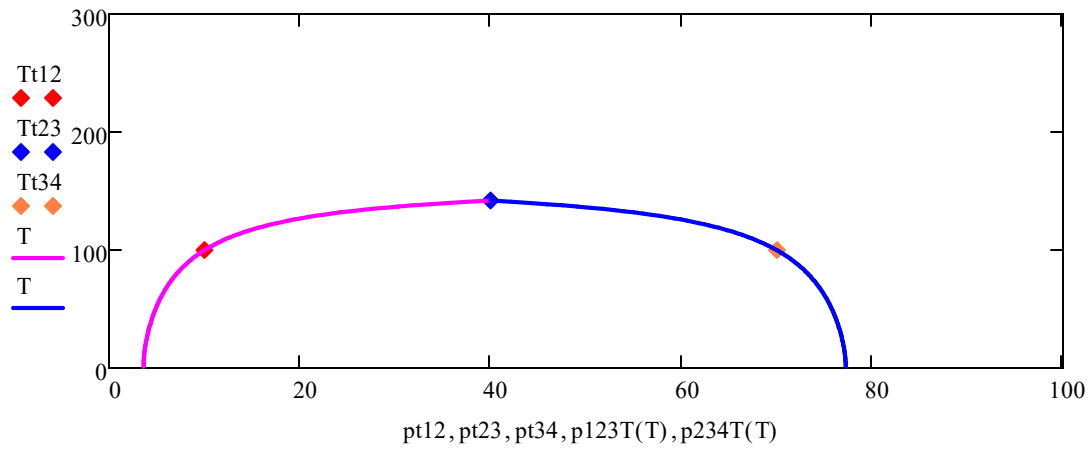
$T := 1, 2 \dots Tt23$



$fv12 := .0063$ $fv34 := .0062$

$$p123T(T) := pt12 + (pt23 - pt12) \cdot \frac{FET(T, Tnm, fv12) - FET(Tt12, Tnm, fv12)}{FET(Tt23, Tnm, fv12) - FET(Tt12, Tnm, fv12)}$$

$$p234T(T) := pt34 + (pt23 - pt34) \cdot \frac{FET(T, Tnm, fv34) - FET(Tt34, Tnm, fv34)}{FET(Tt23, Tnm, fv34) - FET(Tt34, Tnm, fv34)}$$



Kechin-form: $TmKo(p) := Tmo \cdot (1 + qa \cdot p)^{qb}$ $Tmo := 100$ $qa := .5$ $qb := .5$ $qt := .1$
 $p := 0, .1.. 10$

Modified for phase transitions: $TmK1(p, q1) := Tmo \cdot (1 + qa \cdot p)^{qb} \cdot (1 - p \cdot qt)^{q1}$
 $q11 := .18$ $q12 := .1$ $q21 := .25$ $q22 := .5$ $TmK2(p, q2) := Tmo \cdot [1 + qa \cdot p \cdot (1 - p \cdot qt)^2]^{qb}$

